

CALCULATION OF VISCOSITY OF GAS MIXTURES AT ATMOSPHERIC PRESSURE

P. M. Kessel'man and A. S. Litvinov

Inzhenerno-Fizicheskii Zhurnal, Vol. 10, No. 3, pp. 385-392, 1966

UDC 533.16

A method is presented for calculating the viscosity of a multicomponent mixture, based on the use of potential functions with variable parameters $\sigma(T)$ and $\epsilon(T)$. It is shown that the agreement of theory with experiment is good.

Since the available experimental data on viscosity of gas mixtures is limited, theoretical methods which would allow the viscosity η_{ml} to be determined from data on the pure constituents are very important, especially at high temperatures. On the basis of the strict kinetic theory of multicomponent gas mixtures the viscosity of a ν -component mixture is given by the expression [1]

$$\eta_{ml} = \frac{\begin{vmatrix} H_{11} & H_{12} & \dots & H_{1\nu} & x_1 \\ H_{12} & H_{22} & \dots & H_{2\nu} & x_2 \\ \vdots & \vdots & & \vdots & \vdots \\ H_{1\nu} & H_{2\nu} & \dots & H_{\nu\nu} & x_\nu \\ x_1 & x_2 & \dots & x_\nu & 0 \end{vmatrix}}{\begin{vmatrix} H_{11} & H_{12} & \dots & H_{1\nu} \\ H_{12} & H_{22} & \dots & H_{2\nu} \\ \vdots & \vdots & & \vdots \\ H_{1\nu} & H_{2\nu} & \dots & H_{\nu\nu} \end{vmatrix}}, \quad (1)$$

where

$$H_{ii} = \frac{x_i^2}{\eta_i} + \sum_{k=1, k \neq i}^{\nu} \frac{2x_i x_k}{\eta_{ik}} \frac{M_i M_k}{(M_i + M_k)^2} \left[\frac{5}{3A_{ik}^*} + \frac{M_k}{M_i} \right],$$

$$H_{ij} = - \frac{2x_i x_j}{\eta_{ij}} \frac{M_i M_j}{(M_i + M_j)^2} \left[\frac{5}{3A_{ij}^*} - 1 \right], \quad i \neq j.$$

The viscosities of pure substances η_i and the quantities η_{ij} are given by the expressions

$$\eta_i \cdot 10^7 = 266.93 \frac{\sqrt{T M_i}}{\sigma_i^2 \Omega_i^{(2,2)*} (T_i^*)} f_i^{(3)}, \quad (2)$$

$$\eta_{ij} \cdot 10^7 = 266.93 \frac{\sqrt{2M_i M_j T / (M_i + M_j)}}{\sigma_{ij}^2 \Omega_{ij}^{(2,2)*} (T_{ij}^*)} f_{ij}^{(3)}. \quad (3)$$

In these formulas, x_i and M_i are the mole fraction and molecular weight of the i -th component; T is the absolute temperature; $T^* = kT/\epsilon$ is the reduced temperature; σ and ϵ are the parameters of the intermolecular interaction function; $A^* = \Omega^{(2,2)*} / \Omega^{(1,1)*}$ is the ratio of relative collision integrals.

In spite of a certain awkwardness, formula (1) has the undoubted advantage of being based on theory. It may therefore be used for any temperature region (including also high temperatures), and for any mixture, regardless of the number of components. It should be noted that (1) gives a first approximation for the viscosity of a mixture. However, as has been shown in [2], the difference from the second approximation is usually a quantity on the order of 1%. It may therefore be neglected without detriment to the accuracy of the calculations.

The use of (1)-(3) is nevertheless restricted because, for mixtures with components widely different in molecular structure, there is no single potential which would permit correct computation of the interaction between the dissimilar molecules of the components. In addition, in many cases, with complex substances, it is not possible to choose one or the other of the known interaction potentials and thence calculate the viscosity η_i .

The object of the present paper is to show that these difficulties may be overcome by using a method developed by one of the authors [3], which allows a uniform description of all the thermophysical properties of a substance. This method is based on the possibility of representing the intermolecular interaction of any complex molecules by some mean potential function with temperature-dependent potential parameters:

$$u = 4\epsilon(T) \{ [\sigma(T)/r]^{12} - [\sigma(T)/r]^6 \}. \quad (4)$$

In the first place this procedure successfully describes the properties of the individual substances (components) at moderate and high temperatures with an accuracy sufficient for practical purposes. The method has shown that the single potential parameters $\sigma(T)$ and $\epsilon(T)$ give a reliable description not only of the equilibrium properties, but also of the transport coefficients of the gases over a wide range of parameters, including the high-temperature region. This has been verified for a large number of different substances (about 30), composed of both simple and complex molecules, including the polar gases H_2O , NH_3 , HF , HCl , and others. Thus, in calculating the viscosity of a mixture consisting of components of different complexity, it is possible to obtain reliable data on the viscosity η_i over a wide enough temperature range.

Secondly, the fact that any intermolecular interaction, independently of the structure of the colliding

Table 1
Comparison of Calculated Values of Viscosity of H₂-CO₂
Mixtures with Experimental Values [4]

| x _{H₂} , % | η · 10 ⁷ , g/cm · sec | | | | | |
|--------------------------------|----------------------------------|-------------|------------|-------------|------------|-------------|
| | T=400 °K | | T=500 °K | | T=550 °K | |
| | experiment | calculation | experiment | calculation | experiment | calculation |
| 0 | 1944 | 1897 | 2353 | 2312 | 2556 | 2515 |
| 11.79 | 1951 | 1899 | 2360 | 2308 | 2554 | 2508 |
| 19.93 | 1945 | 1897 | 2358 | 2302 | 2542 | 2500 |
| 41.29 | 1933 | 1878 | 2321 | 2265 | 2506 | 2452 |
| 59.46 | 1878 | 1824 | 2239 | 2184 | 2471 | 2356 |
| 78.50 | 1713 | 1669 | 2026 | 1976 | 2173 | 2122 |
| 88.88 | 1526 | 1483 | 1783 | 1740 | 1904 | 1861 |
| 100.00 | 1081 | 1090 | 1256 | 1264 | 1341 | 1346 |
| mean error, % | | 2.44 | | 2.09 | | 2.10 |

molecules, is described by the single potential function (4), gives a basis for the use of the function to describe the interaction of nonhomogeneous molecules. This makes it possible to calculate coefficients in accordance with (3), in which the collision integrals $\Omega^{(2,2)*}$, tabulated in [1] for the potential (4), may be used. For this purpose it is necessary to know the potential parameters σ_{ij} and ϵ_{ij} . When components *i* and *j* are composed of spherical molecules, σ_{ij} and ϵ_{ij} , as is known, may be reliably determined according to the following empirical relations:

$$\sigma_{ij} = \frac{1}{2} (\sigma_i + \sigma_j), \quad \epsilon_{ij} = \sqrt{\epsilon_i \epsilon_j}. \quad (5)$$

Since the model corresponding to potential (4) with variable parameters $\sigma(T)$ and $\epsilon(T)$ is a generalized model of a pseudo-spherical gas (with centrally-symmetric forces), for a mixture with any components the most probable rules for combining the potential parameters must be the rules (5).

In the present work we calculated the viscosity of mixtures according to (1)-(3), using the intermolecular interaction potential functions (4) and the combination rules (5). For this purpose it was necessary to know only the potential parameters, $\sigma(T)$ and $\epsilon(T)$, of the pure components. These parameters are given in Tables 3 and 4 for some gases. Viscosities were calculated for 20 binary mixtures and for a number of multicomponent mixtures for which experimental data appear in the literature, and a comparison with experiment was made. The mixtures comprised gases of different complexities, H₂, O₂, N₂, CO₂, CH₄, C₂H₄, NH₃, H₂O, etc.

It should be noted that for all the mixtures checked, both binary and multicomponent (Fig. 4), good agreement (3% or better) with experiment was found. It is important to stress here that good results were also obtained for systems which had a far from linear dependence of viscosity on composition (H₂-CO₂, Table 1).

There is naturally very great interest in the results of comparison with mixtures whose components correspond to various molecular models and cannot be described by the classical Lennard-Jones potentials with constants ϵ and σ .

In this respect mixtures containing a polar component are typical. For all such complex systems on which there are experimental data (NH₃-H₂, NH₃-O₂, NH₃-N₂, NH₃-C₂H₄, etc.), good agreement with experiment (1-3%) is observed. As an illustrative example we shall make a comparison with the experimental data of [4] for the mixture NH₃-C₂H₄ in Table 2 and NH₃-N₂ in Fig. 1.

An exception is the unique mixture of 50% H₂O and 50% N₂ shown in Fig. 2, for which considerable discrepancies (on the order of 15-20%) with the experimental data [5] were observed. As regards the data of [5], it should be noted that they fall outside the general relation observed for the viscosity of binary mixtures. In the first place, the fact should be borne in mind that these data practically coincide with the viscosity of pure water vapor. Secondly, if, using the data of [5], we try to construct the dependence of mixture viscosity on composition, as in Fig. 3, the curve $\eta_m = f(x)$ is seen to be anomalous (there is a steep

Table 2
Comparison of Calculated Values of Viscosity of
NH₃-C₂H₄ Mixtures with Experimental Data [4]

| x _{C₂H₄} , % | η · 10 ⁷ , g/cm · sec | | | |
|---|----------------------------------|-------------|-------------|-------------|
| | T=473.15° K | | T=523.15° K | |
| | experiment | calculation | experiment | calculation |
| 0 | 1646 | 1646 | 1813 | 1848 |
| 11.33 | 1647 | 1631 | 1809 | 1821 |
| 19.29 | 1648 | 1621 | 1805 | 1805 |
| 30.39 | 1639 | 1608 | 1791 | 1783 |
| 42.28 | 1622 | 1590 | 1764 | 1750 |
| 70.07 | 1595 | 1571 | 1729 | 1713 |
| 89.04 | 1561 | 1558 | 1689 | 1685 |
| 100.00 | 1541 | 1550 | 1666 | 1670 |
| mean error, % | | 1.09 | | 0.66 |

Table 3
Potential Parameters ϵ/k , °K and σ , Å

| T, °K | For H ₂ O | | For NH ₃ | | For CO ₂ | |
|-------|----------------------|----------|---------------------|----------|---------------------|----------|
| | ϵ/k | σ | ϵ/k | σ | ϵ/k | σ |
| 300 | — | — | — | — | 222.1 | 3.931 |
| 350 | — | — | — | — | 220.5 | 3.927 |
| 400 | 632.6 | 2.971 | 377.9 | 3.251 | 219.0 | 3.925 |
| 450 | 620.2 | 2.958 | 366.9 | 3.236 | 217.4 | 3.920 |
| 500 | 609.0 | 2.946 | 356.4 | 3.221 | 215.7 | 3.916 |
| 550 | 599.1 | 2.936 | 346.7 | 3.207 | 213.8 | 3.913 |
| 600 | 589.9 | 2.926 | 337.3 | 3.194 | 211.3 | 3.910 |
| 650 | 580.6 | 2.918 | 327.2 | 3.182 | 208.4 | 3.908 |
| 700 | 572.0 | 2.911 | 315.2 | 3.171 | 206.4 | 3.906 |
| 750 | 564.0 | 2.904 | 303.2 | 3.161 | 206.9 | 3.904 |
| 800 | 556.3 | 2.898 | 294.0 | 3.152 | 208.0 | 3.902 |
| 850 | 548.5 | 2.892 | 287.6 | 3.144 | 209.4 | 3.901 |
| 900 | 541.0 | 2.886 | 283.6 | 3.137 | 210.5 | 3.899 |
| 950 | 532.7 | 2.881 | 282.0 | 3.131 | 211.4 | 3.898 |
| 1000 | 523.5 | 2.876 | 282.1 | 3.125 | 212.1 | 3.897 |
| 1100 | 502.5 | 2.868 | 283.1 | 3.115 | 213.1 | 3.894 |
| 1200 | 481.1 | 2.860 | 285.1 | 3.106 | 213.8 | 3.891 |
| 1300 | 466.1 | 2.852 | 287.7 | 3.097 | 214.3 | 3.889 |
| 1400 | 458.0 | 2.845 | 290.6 | 3.090 | 214.7 | 3.887 |
| 1500 | 454.1 | 2.839 | 293.5 | 3.083 | 215.0 | 3.885 |
| 1600 | 453.5 | 2.834 | 296.2 | 3.076 | 215.2 | 3.883 |
| 1700 | 454.4 | 2.828 | 298.5 | 3.070 | 215.5 | 3.881 |
| 1800 | 456.0 | 2.824 | 300.5 | 3.065 | 215.7 | 3.879 |
| 1900 | 458.1 | 2.819 | 302.2 | 3.060 | 215.9 | 3.877 |
| 2000 | 461.0 | 2.815 | 303.6 | 3.055 | 216.1 | 3.875 |
| 2200 | 467.4 | 2.807 | 306.0 | 3.045 | 216.3 | 3.872 |
| 2400 | 473.2 | 2.800 | 307.8 | 3.037 | 216.5 | 3.869 |
| 2600 | 478.2 | 2.794 | 309.2 | 3.029 | 216.7 | 3.867 |
| 2800 | 482.2 | 2.788 | 310.2 | 3.022 | 216.8 | 3.865 |
| 3000 | 485.5 | 2.782 | 311.2 | 3.016 | 216.9 | 3.863 |

| T, °K | For C ₂ H ₄ | | For N ₂ | | For O ₂ | |
|-------|-----------------------------------|----------|--------------------|----------|--------------------|----------|
| | ϵ/k | σ | ϵ/k | σ | ϵ/k | σ |
| 200 | — | — | 95.89 | 3.702 | 117.8 | 3.467 |
| 250 | — | — | 95.57 | 3.702 | 117.6 | 3.467 |
| 300 | 213.9 | 4.235 | 95.16 | 3.701 | 117.3 | 3.466 |
| 350 | 212.6 | 4.231 | 95.01 | 3.701 | 116.9 | 3.466 |
| 400 | 211.4 | 4.227 | 95.25 | 3.700 | 116.5 | 3.465 |
| 450 | 210.2 | 4.223 | 95.41 | 3.700 | 116.7 | 3.465 |
| 500 | 208.9 | 4.220 | 95.51 | 3.699 | 116.9 | 3.464 |
| 550 | 207.4 | 4.217 | 95.59 | 3.699 | 117.0 | 3.464 |
| 600 | 205.4 | 4.214 | 95.64 | 3.699 | 117.1 | 3.464 |
| 650 | 203.2 | 4.212 | 95.68 | 3.699 | 117.2 | 3.464 |
| 700 | 202.6 | 4.210 | 95.71 | 3.699 | 117.3 | 3.463 |
| 750 | 202.9 | 4.209 | 95.72 | 3.699 | 117.3 | 3.463 |
| 800 | 203.9 | 4.208 | 95.75 | 3.699 | 117.4 | 3.463 |
| 850 | 205.0 | 4.206 | 95.74 | 3.698 | 117.4 | 3.463 |
| 900 | 205.9 | 4.205 | 95.74 | 3.698 | 117.4 | 3.463 |
| 950 | 206.6 | 4.204 | 95.75 | 3.698 | 117.4 | 3.462 |
| 1000 | 207.2 | 4.203 | 95.75 | 3.698 | 117.4 | 3.462 |
| 1100 | 208.0 | 4.202 | 95.75 | 3.698 | 117.4 | 3.462 |
| 1200 | 208.4 | 4.201 | 95.76 | 3.698 | 117.5 | 3.462 |
| 1300 | 208.7 | 4.200 | 95.76 | 3.698 | 117.5 | 3.462 |
| 1400 | 208.9 | 4.199 | 95.76 | 3.697 | 117.5 | 3.462 |
| 1500 | 209.0 | 4.198 | 95.76 | 3.697 | 117.5 | 3.462 |
| 1600 | 209.1 | 4.197 | 95.76 | 3.697 | 117.5 | 3.461 |
| 1700 | 209.2 | 4.196 | 95.76 | 3.697 | 117.5 | 3.461 |
| 1800 | 209.3 | 4.195 | 95.76 | 3.697 | 117.5 | 3.461 |
| 1900 | 209.4 | 4.195 | 95.76 | 3.697 | 117.5 | 3.461 |
| 2000 | 209.4 | 4.194 | 95.76 | 3.697 | 117.5 | 3.461 |
| 2200 | 209.5 | 4.192 | 95.76 | 3.697 | 117.5 | 3.461 |
| 2400 | 209.5 | 4.191 | 95.76 | 3.696 | 117.5 | 3.461 |
| 2600 | 209.5 | 4.189 | 95.76 | 3.696 | 117.5 | 3.460 |
| 2800 | 209.5 | 4.188 | 95.76 | 3.696 | 117.5 | 3.460 |
| 3000 | 209.5 | 4.187 | 95.76 | 3.696 | 117.5 | 3.460 |

Table 4
Potential Parameters at the Boyle Temperature

| Gas | T_B , °K | $\frac{a_B}{k}$, °K | σ_B , Å |
|-------------------------------|------------|----------------------|----------------|
| H ₂ O | 1550 | 453.6 | 2.836 |
| NH ₃ | 963.3 | 282.0 | 3.129 |
| CO ₂ | 705.0 | 206.4 | 3.905 |
| C ₂ H ₄ | 692.1 | 202.6 | 4.211 |
| N ₂ | 324.4 | 94.96 | 3.701 |
| O ₂ | 397.9 | 116.5 | 3.465 |

Table 5
Viscosity of the Pure Components

| T , °K | $\eta \cdot 10^7$, g/cm · sec | | | | | |
|----------|--------------------------------|-----------------|-----------------|-------------------------------|----------------|----------------|
| | H ₂ O | NH ₃ | CO ₂ | C ₂ H ₄ | N ₂ | O ₂ |
| 200 | — | — | — | — | 1260 | 1424 |
| 250 | — | — | — | — | 1515 | 1728 |
| 300 | — | — | 1444 | 1020 | 1748 | 2010 |
| 350 | — | — | 1675 | 1173 | 1962 | 2272 |
| 400 | 1277 | 1352 | 1897 | 1321 | 2160 | 2513 |
| 450 | 1466 | 1555 | 2108 | 1468 | 2348 | 2742 |
| 500 | 1661 | 1757 | 2312 | 1608 | 2527 | 2957 |
| 550 | 1858 | 1959 | 2515 | 1740 | 2698 | 3163 |
| 600 | 2057 | 2157 | 2701 | 1869 | 2861 | 3359 |
| 650 | 2258 | 2358 | 2887 | 1992 | 3021 | 3551 |
| 700 | 2458 | 2559 | 3061 | 2110 | 3175 | 3736 |
| 750 | 2657 | 2756 | 3223 | 2224 | 3325 | 3916 |
| 800 | 2853 | 2936 | 3378 | 2333 | 3470 | 4090 |
| 850 | 3048 | 3126 | 3526 | 2436 | 3612 | 4260 |
| 900 | 3242 | 3292 | 3673 | 2536 | 3751 | 4426 |
| 950 | 3436 | 3445 | 3815 | 2636 | 3881 | 4588 |
| 1000 | 3630 | 3593 | 3955 | 2733 | 4020 | 4746 |
| 1100 | 4015 | 3879 | 4229 | 2979 | 4277 | 5054 |
| 1200 | 4392 | 4144 | 4494 | 3098 | 4529 | 5353 |
| 1300 | 4754 | 4404 | 4750 | 3272 | 4771 | 5642 |
| 1400 | 5084 | 4652 | 4998 | 3440 | 5006 | 5920 |
| 1500 | 5391 | 4895 | 5241 | 3603 | 5235 | 6193 |
| 1600 | 5680 | 5129 | 5476 | 3762 | 5460 | 6458 |
| 1700 | 5958 | 5359 | 5706 | 3918 | 5680 | 6717 |
| 1800 | 6228 | 5587 | 5931 | 4070 | 5895 | 6972 |
| 1900 | 6488 | 5809 | 6152 | 4217 | 6103 | 7221 |
| 2000 | 6742 | 6028 | 6369 | 4364 | 6311 | 7466 |
| 2200 | 7230 | 6458 | 6790 | 4648 | 6710 | 7941 |
| 2400 | 7704 | 6871 | 7198 | 4925 | 7102 | 8403 |
| 2600 | 8162 | 7277 | 7590 | 5193 | 7478 | 8849 |
| 2800 | 8610 | 7676 | 7973 | 5450 | 7846 | 9285 |
| 3000 | 9048 | 8060 | 8346 | 5703 | 8204 | 9709 |

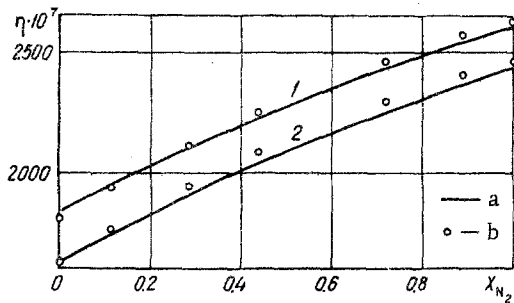


Fig. 1. Dependence of the viscosity of the mixture NH_3-N_2 (g/cm·sec) on composition and temperature: 1) at 523.15° K; 2) 473.15° K; a) data of present paper; b) experimental data of [4].

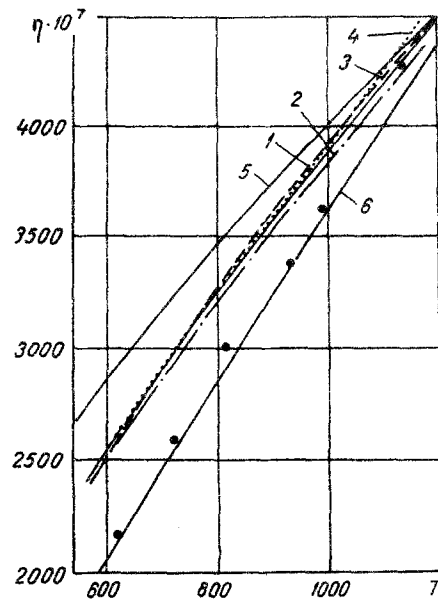


Fig. 2. Dependence on T (°K) of the viscosity η (g/cm·sec) of a mixture of 50% (molar) H_2O and 50% N_2 : 1, 2, 3, 4) calculated curves from the data of the present paper, from the Wilke formula [6], and by the methods of [7] and [8], respectively; 5, 6) experimental curves for N_2 and H_2O ; the points are experimental of [5].

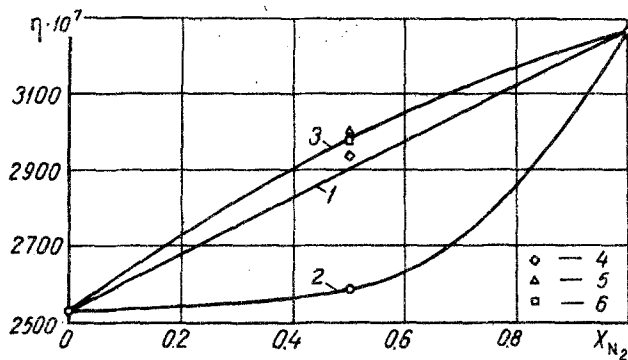


Fig. 3. Dependence of viscosity η (g/cm·sec) of a 50% H_2O -50% N_2 mixture on composition at temperature 718.55° K: 1) according to a linear law; 2) according to the experimental data of [5]; 3) according to the data of the present paper; 4) according to the Wilke formula [6]; 5) according to the method of reference [7]; 6) according to the method of [8].

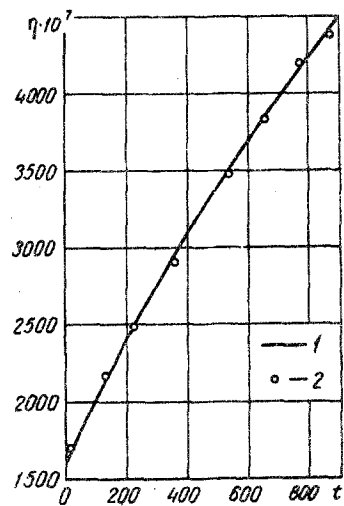


Fig. 4. Dependence of viscosity η (g/cm·sec) of a mixture of 19.82% CO_2 -79.54% N_2 -0.64% O_2 on temperature t (°C): 1) according to the data of the present paper; 2) according to the experimental data of [9].

negative slope of η_m with viscosity, calculated according to a linear law). For binary mixtures, as a rule, the observed behavior of the curve is the opposite (Fig. 1).

Figure 2 also gives values of viscosity calculated by other independent methods, described in references [6, 7, 8], and our data practically coincide with the results of these calculations, the discrepancies being not more than 2%. This in turn confirms the correctness of our calculations, and is also evidence that the data of [5] should be considered appreciably too low.

Tables 3 and 4 give the potential parameters of several gases used for calculating the viscosities of the mixtures illustrated in Figs. 1-4. These parameters were found by the method described in [3]. It may be seen from the table, that the dependence of the potential parameters on temperature is the stronger, the more complex the gas molecule. For diatomic gases (N_2, O_2) the potential parameters do not change appreciably with change of temperature, since their molecules do not possess appreciable asymmetry. Constant potential parameters of these gases have been given repeatedly in the literature, but they were suitable for calculating either the second virial coefficient alone, or the viscosity alone. Although the σ and ϵ/k given in the present paper for diatomic gases may be considered practically constant for calculating viscosity, their value rests in the fact that they are unique for calculating both transport and thermodynamic properties. Moreover, it should be kept in mind that the quantity σ enters into the formulas for calculating the second and third virial coefficients at the third and sixth power, respectively. Therefore a small error in σ may lead to large errors in calculating the virial coefficients. Table 5 gives the viscosities of pure gases calculated on the basis of the potential parameters obtained in this paper.

Thus, the following conclusions may be drawn:

1. The use of the potential function (4) with variable potential parameters permits a reliable description of the viscosity not only of pure components, but also of mixtures.

2. The method of calculating the viscosity of a mixture, using the exact formulas of kinetic theory, the proposed potential function with variable parameters, and the combination rules (5), is a universal one, gives high accuracy in practice, and may be used for any multicomponent mixture over a wide temperature range, including the high-temperature region.

REFERENCES

1. J. Hirschfelder, C. Curtiss, and R. Bird, *Molecular Theory of Gases and Liquids* [Russian translation], IL, 1961.
2. S. C. Saxena and R. K. Joshi, *Indian J. Phys.*, **37**, 9, 479, 1963.
3. P. M. Kessel'man, *Teplofizika vysokikh temperatur*, **2**, no. 6, 1964.
4. M. Trautz and K. Sorg, *Ann. Phys.*, **10**, 81, 1931.
5. C. F. Bonilla, S. E. Greer and E. A. Taikoff, *Thermodynamic and Transport Properties of Gases Liquids and Solids*. McGraw-Hill, New York, 346-349, 1959.
6. G. R. Wilke, *J. Chem. Phys.*, **18**, 517, 1950.
7. H. Mann, *Gas-u-Wasserfach*, **73**, 570, 1930.
8. C. F. Bonilla, R. D. Brooks and P. L. Walker, *Proceedings of the General Discussion on Heat Transfer*. The Institution of Mechanical Engineers, London, 163-173, 1951.
9. M. J. Kenney, R. J. Sarjant and M. W. Thring, *J. Appl. Phys.*, **7**, no. 9, 324, 1956.

3 July 1965

Odessa Lomonosov Technological Institute